

2.2 Fokker-Planck-Gleichung

Zeitentwicklung eines kontinuierlichen Markov-Prozesses:
 $X(t)$ (1-dim. Zufallsvar.)

$$\frac{\partial}{\partial t} p(x, t | x_0, t_0) = \underbrace{-\frac{\partial}{\partial x} [A(x, t) p(x, t | x_0, t_0)]}_{\text{Drift}} + \underbrace{\frac{1}{2} \frac{\partial^2}{\partial x^2} [B(x, t) p(x, t | x_0, t_0)]}_{\text{Diff.}}$$

Anfangsbed. $p(x, t_0 | x_0, t_0) = \delta(x - x_0)$

$\hat{=}$ Ein-Zeit-Wahrsch. $p(x, t) = \int dx_0 p(x, t | x_0, t_0) = \int dx_0 p(x, t | x_0, t_0) p(x_0, t_0)$
 mit Anf.-bed. $p(x, t_0)$ (weniger singular)

Randbed. (n-dim.)

Fokker-Planck (FP)-gl. ist lokale Bilanzgl.

$$\frac{\partial p(x, t)}{\partial t} + \sum_i \frac{\partial}{\partial x_i} J_i(x, t) \quad (\dot{p} + \text{div} \underline{J} = 0)$$

mit Wahrscheinl.strom

$$\underline{J}_i(x, t) = A_i(x, t) p(x, t) - \frac{1}{2} \sum_j \frac{\partial}{\partial x_j} (B_{ij}(x, t) p(x, t))$$

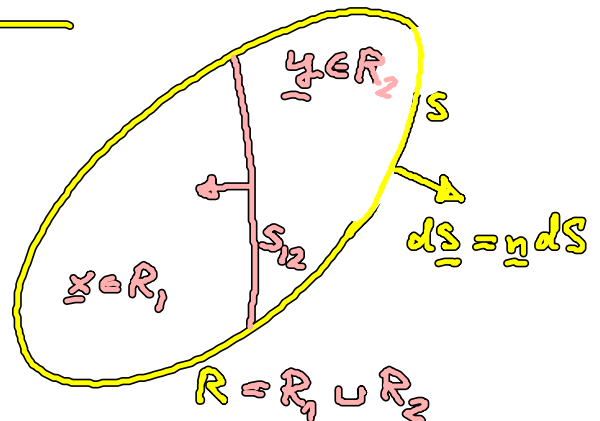
(vgl. Teilchendichte)

$$\begin{cases} \dot{\rho} + \text{div} \underline{J} = 0 \\ \underline{J} = \underline{v} \rho - D \nabla \rho \\ \Rightarrow \dot{\rho} = -\rho \nabla \cdot \underline{v} + D \Delta \rho \end{cases}$$

Globale Bilanzgl. für Gebiet $R \in \mathbb{R}^n$

$$P(R, t) := \int_R dx p(x, t)$$

$$\frac{\partial P}{\partial t} = - \int_S d\underline{S} \cdot \underline{J}(x, t) \quad (\text{Gauss'scher Satz})$$



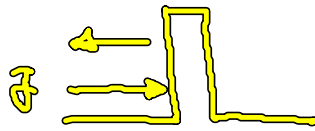
Netto-Wahrsch. flux durch beliebige Fläche S_{12} :

$$\lim_{\Delta t \rightarrow 0} \frac{1}{\Delta t} \int_{R_1} dx \int_{R_2} dy [p(x, t + \Delta t, y, t) - p(y, t + \Delta t, x, t)] = \int_{S_{12}} dS \cdot \underline{j}(z, t)$$

$R_1 \leftarrow R_2$ S_{12}

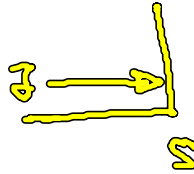
(a) Reflektierende Barriere:

$$\underline{n} \cdot \underline{j}(x, t) \Big|_{x \in S} = 0$$



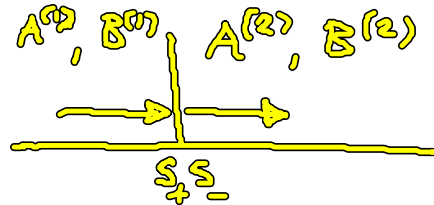
(b) Absorbierende Barriere:

$$p(x, t) \Big|_{x \in S} = 0$$



(c) Grenzfläche zwischen 2 Medien: $A^{(1)}, B^{(1)}$ $A^{(2)}, B^{(2)}$

$$\underline{n} \cdot \underline{j} \Big|_{s_+} = \underline{n} \cdot \underline{j} \Big|_{s_-}$$

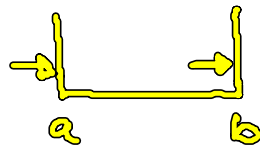


$$p \Big|_{s_+} = p \Big|_{s_-}$$

(d) Periodische Randbed.:

$$p(a, t) = p(b, t)$$

$$j(a, t) = j(b, t)$$



(e) Natürliche Randbed.: $A(a, t) = 0$

(geschw. = 0, Rand wird nie erreicht)

Stationäre Lösung für homogenen Markov-Prozess

homogen $\Rightarrow A, B$ unabh. von t

1-dim.: $\frac{d}{dx} j(x, t) = 0 \Rightarrow j(x) = \text{const.} = j(a) = j(b)$

(i) reflekt. Randbed. $\Rightarrow j(x) = 0$

$$\Rightarrow A(x)p^*(x) = \frac{1}{2} \frac{d}{dx} [B(x)p^*(x)] = 0 \quad \Rightarrow 2 \frac{A}{B} dx = \frac{d(Bp^*)}{Bp^*}$$

$$\Rightarrow p^*(x) = \frac{N}{B(x)} \exp \left[2 \int_a^x \frac{A(x')}{B(x')} dx' \right]$$

Intenziallösung, Normierung $\int_a^b dx p^*(x) = 1 \Rightarrow N$

(ii) period. Randbed. $\Rightarrow J(x) = J$

$$A(x)p^*(x) - \frac{1}{2} \frac{d}{dx} [B(x)p^*(x)] = J \quad (1) \text{ lin. inhom. Dgl.}$$

Mit $\psi(x) := \exp \left[2 \int_a^x \frac{A(x')}{B(x')} dx' \right]$ ergibt sich

$$A \frac{\psi}{B} - \frac{1}{2} \frac{d}{dx} \psi = 0 \quad (\text{homog. L\u00f6s } p^* = \frac{\psi}{B})$$

$$\Rightarrow A = \frac{B}{2} \frac{\psi'}{\psi} \stackrel{in(1)}{\Rightarrow} B p^* \frac{\psi'}{\psi} - (B p^*)' = 2J$$

$$\Rightarrow \frac{-(B p^*) \psi' + (B p^*)' \psi}{\psi^2} = -\frac{2J}{\psi}$$

$$\Rightarrow \left(\frac{B p^*}{\psi} \right)' = -\frac{2J}{\psi}$$

$$\int_a^x dx' \frac{d}{dx'} \left(\frac{B p^*}{\psi} \right) = -2J \int_a^x \frac{dx'}{\psi(x')} \Rightarrow p^*(x)$$

Bestimmung von J durch period. Randbed.

$$\Rightarrow p^*(x) = p^*(a) \frac{\int_a^x \frac{dx'}{\psi(x')} \frac{B(b)}{\psi(b)} + \int_x^b \frac{dx'}{\psi(x')} \frac{B(a)}{\psi(a)}}{\frac{B(x)}{\psi(x)} \int_a^b \frac{dx'}{\psi(x')}} \quad \text{[Note: The original image has a typo in the denominator of the fraction above, it should be } \int_a^b \frac{dx'}{\psi(x')} \text{]} \Rightarrow p^*(x) = p^*(a) \frac{\int_a^x \frac{dx'}{\psi(x')} \frac{B(b)}{\psi(b)} + \int_x^b \frac{dx'}{\psi(x')} \frac{B(a)}{\psi(a)}}{\int_a^b \frac{dx'}{\psi(x')} \frac{B(x)}{\psi(x)}}$$

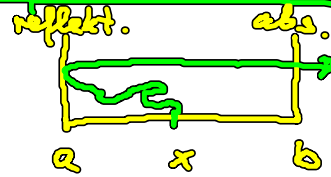
First passage time

Fragestellung: Wie lange hält sich ein Teilchen in einem vorgegebenen Gebiet auf?



Teilchen zwischen 1 absorb. u. 1 reflektierenden Barriere

Ziel: Entweichzeit T (escape time)



Wahrsch., dass das Teilchen zur Zeit t noch in (a, b) , wenn es bei x gestartet ist:

$$G(x, t) := \int_a^b dx' p(x', t | x, 0) \equiv \text{Prob}(T \geq t)$$

stationärer Prozess: $p(x', t | x, 0) = p(x', 0 | x, -t)$

Rückwärts-FP-gle.: Rückwärtsentw. für $t' \leq t$ aus (x, t)

$$\frac{\partial p(x, t | y, t')}{\partial t'} = -A(y, t') \frac{\partial p(x, t | y, t')}{\partial y} - \frac{1}{2} B(y, t') \frac{\partial^2 p(x, t | y, t')}{\partial y^2}$$

homog. (A, B zeitunabh.):
$$\frac{\partial}{\partial t} G(x, t) = - \frac{\partial}{\partial t'} \int_a^b dx' p(x', 0 | x, t')$$

$$\frac{\partial}{\partial t} G(x, t) = A(x) \frac{\partial}{\partial x} G(x, t) + \frac{1}{2} B(x) \frac{\partial^2}{\partial x^2} G(x, t)$$

„End“ bed.: $p(x', 0 | x, 0) = \delta(x - x')$

$$\Rightarrow G(x, 0) = \begin{cases} 1 & a \leq x \leq b \\ 0 & \text{sonst} \end{cases}$$

absorb. Randbed.: $\text{Prob}(T \geq t) = 0$ wenn $x = a$ oder b

$$\Leftrightarrow G(a, t) = G(b, t) = 0$$

mittlere erste Übergangszeit (mean first passage time)

$$T(x) := \langle T \rangle = - \int_0^{\infty} t dG = - \int_0^{\infty} dt t \frac{\partial}{\partial t} G(x, t) \stackrel{\text{part. int.}}{=} \int_0^{\infty} dt G(x, t)$$

Bgl. für $T(x)$ aus der Rückwärts-TP-Gl.:

$$\int_0^{\infty} \frac{\partial}{\partial t} G(x,t) dt = \underbrace{G(x,\infty)}_0 - \underbrace{G(x,0)}_1 = -1$$

** $A(x) \frac{\partial}{\partial x} T(x) + \frac{1}{2} B(x) \frac{\partial^2}{\partial x^2} T(x) = -1$, Randbed. $T(a) = T(b) = 0$

Lösung der hom.-Gl. $A\psi - \frac{1}{2} B \frac{\partial}{\partial x} \psi = 0 \Leftrightarrow 2 \frac{A}{B} dx = \frac{d\psi}{\psi}$

$$\Leftrightarrow \psi(x) = \exp \left[2 \int_a^x dx' \frac{A(x')}{B(x')} \right]$$

\Rightarrow Lösung der inhom. Gl. $T(x)$ durch $\psi(x)$ ausgedrückt:

$$T(x) = 2 \int_x^b \frac{dy}{\psi(y)} \int_a^y dz \frac{\psi(z)}{B(z)} \quad (*)$$

Beweis: $T' = -\frac{2}{\psi(x)} \int_a^x dz \frac{\psi(z)}{B(z)}$

$$T'' = -\frac{2}{\psi(x)^2} \left[\frac{\psi(x)^2}{B(x)} - \underbrace{\psi'(x)}_a \int_a^x dz \frac{\psi(z)}{B(z)} \right]$$

(**) $\Rightarrow AT' + \frac{1}{2} BT'' = \underbrace{-\frac{2A}{\psi} \int_a^x dz \frac{\psi}{B}}_{-1} - \left[1 - \underbrace{\frac{B}{\psi} \psi' \int_a^x dz \frac{\psi}{B}}_{-1} \right] = -1$

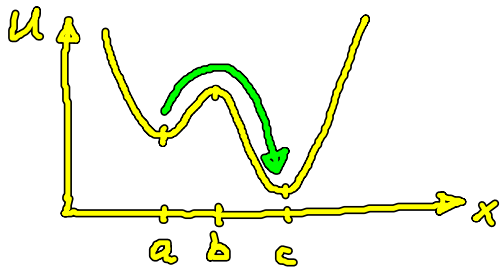
$\psi' = \frac{2A}{B} \psi \Rightarrow \frac{2A}{\psi}$

Randbed.: $T'(a) = \int_0^{\infty} dt \frac{\partial}{\partial x} G(x,t) \Big|_{x=a} = 0 \quad \checkmark$

$T(b) = 0 \quad \checkmark$

□

Anwendung: Kramers' Problem (1940)



Entweichen über Potenzialbarriere:

bistab. Pot. $U(x)$

überdämpftes Teilchen: $\dot{x} = -U'(x) \equiv A(x)$

Diff. konst. $D = \frac{\beta}{2}$

Kraft

$$\text{FP-GE. } \frac{\partial}{\partial t} p(x,t) = \frac{\partial}{\partial x} [U'(x)p(x,t)] + D \frac{\partial^2}{\partial x^2} p(x,t)$$

stat. Lös. $p^*(x) = \sqrt{V} \exp \left[- \int_a^x dx' \frac{U'(x')}{D} \right] = \sqrt{V} \exp \left[- \frac{U(x)}{D} \right]$
(reflekt. Randbed.)