

# English summary

## 2 Phenomenological models

few equations, simple nonlinearities, feasible for bifurcation analysis  
 qualitative agreement of time series

### 2.1 FitzHugh-Nagumo model

$$\epsilon \dot{x} = x - \frac{x^3}{3} - y$$

$x$ : activator

,  $a$ : bifurcation parameter

$$\dot{y} = x + a$$

$y$ : inhibitor

,  $\epsilon$ : timescale separation

$\epsilon \ll 1 \Rightarrow x$  fast,  $y$  slow

$|a| < 1$ : oscillatory (limit cycle)

$|a| > 1$ : excitable (fixed point)

linear  
 $\xrightarrow{\text{stability analysis}}$

fixed point  
 $(\dot{x}=0, \dot{y}=0)$

$$x^* = -a$$

$$y^* = -a + \frac{a^3}{3}$$

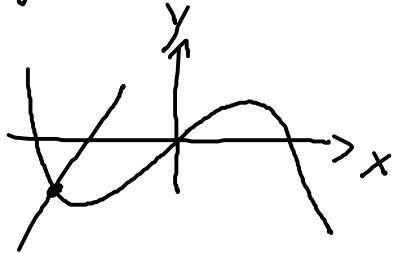
$\rightarrow$  unstable for  $|a| < 1$

$\searrow$  stable for  $|a| > 1$

(extended) FitzHugh-Nagumo model  $\cong$  Buzkoeff-vanderPol model

$$\epsilon \dot{x} = x - \frac{x^3}{3} - y$$

$$\dot{y} = x + a - \gamma y$$



Canard trajectory:

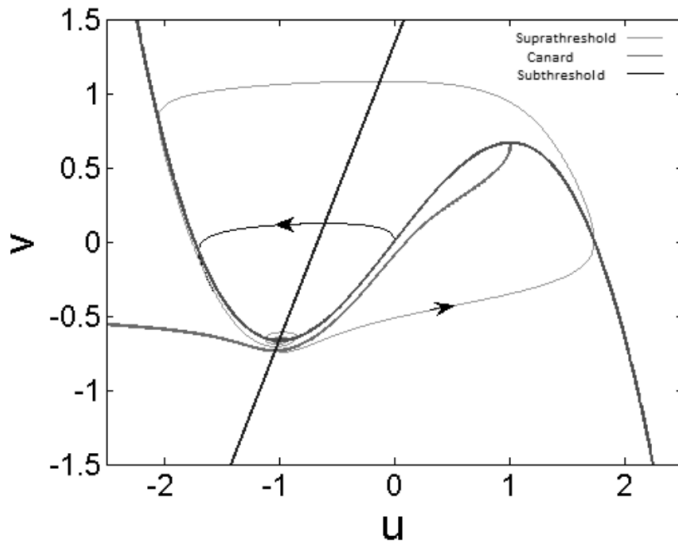
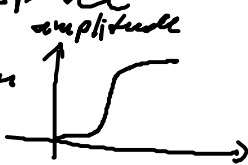
separates subthreshold

oscillations from

large excursions

in phase space

$\Rightarrow$  explosion



excitability type II: at bifurcation point: limit cycle with amplitude 0

and finite frequency  $\text{Im} \lambda \Big|_{|a|=1} = \frac{1}{\sqrt{\epsilon}}$

## 2.2 SNIPER model

Saddle-Node Infinite Period Bifurcation

wird als SNIC, Saddle-Node on an Invariant Cycle bekannt.

$$\left. \begin{aligned} \dot{x} &= x(1-x^2-y^2) + y(x+b) \\ \dot{y} &= y(1-x^2-y^2) - x(x+b) \end{aligned} \right\} \begin{aligned} x &= r \cos \varphi \\ y &= r \sin \varphi \end{aligned} \quad \begin{aligned} \dot{r} &= r(1-r^2) \\ \dot{\varphi} &= b - r \cos \varphi \end{aligned}$$

PHYSICAL REVIEW E

VOLUME 50, NUMBER 5

NOVEMBER 1994

### Resonancelike responses of autonomous nonlinear systems to white noise

T. Ditzinger, C. Z. Ning,\* and G. Hu†

Institut für Theoretische Physik und Synergetik, Universität Stuttgart, Pfaffenwaldring 57/IV,  
D-70550 Stuttgart, Federal Republic of Germany

## PHYSICAL REVIEW LETTERS

VOLUME 71

9 AUGUST 1993

NUMBER 6

### Stochastic Resonance without External Periodic Force

Hu Gang

International Centre for Theoretical Physics, Trieste 34100, Italy  
and Physics Department, Beijing Normal University, Beijing 100875, People's Republic of China

T. Ditzinger,\* C. Z. Ning, and H. Haken

Institut für Theoretische Physik und Synergetik, Universität Stuttgart, Pfaffenwaldring 57/IV,  
D-7000 Stuttgart 80, Federal Republic of Germany  
(Received 21 December 1992)

Bestimmung des Fixpunkts.

$$0 = x(1-x^2-y^2) + y(x-b)$$

$$0 = y(1-x^2-y^2) - x(x-b)$$

trivialer Fixpunkt am Ursprung  $(x_A^*, y_A^*) = (0, 0)$

weitere Fixpunkte für  $x-b=0 \Rightarrow x^*=b$

$$\text{und } 1-x^2-y^2=0 \Rightarrow y^2=1-b^2 \Rightarrow y^* = \pm \sqrt{1-b^2}$$

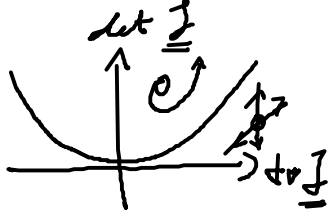
$$\Rightarrow \left. \begin{aligned} (x_B^*, y_B^*) &= (b, \sqrt{1-b^2}) \\ (x_C^*, y_C^*) &= (b, -\sqrt{1-b^2}) \end{aligned} \right\} \begin{aligned} &\text{existieren nur} \\ &\text{für } |b| < 1 \end{aligned} \quad \left. \begin{aligned} (r_B^*, \varphi_B^*) &= (1, \arccos b) \\ (r_C^*, \varphi_C^*) &= (1, -\arccos b) \end{aligned} \right\}$$

Lineare Stabilitätsanalyse:

$$\text{Jacobi-Matrix: } \underline{J} = \begin{pmatrix} 1-3x^2-y^2+y & -2xy+x-b \\ -2xy-2x+b & 1-x^2-3y^2 \end{pmatrix}$$

$$\text{Eigenwerte } \lambda_{1,2} = \frac{\text{tr } \underline{J} \pm \sqrt{(\text{tr } \underline{J})^2 - 4 \det \underline{J}}}{2}$$

$$\text{1. Fall: } (x_A^*, y_A^*) = (0, 0) \Rightarrow \underline{J}_A = \begin{pmatrix} 1 & -b \\ b & 1 \end{pmatrix} \Rightarrow \lambda_{1,2} = \frac{2 \pm \sqrt{4-4(1+b^2)}}{2} = 1 \pm ib$$



$(x_A^*, y_A^*)$  ist ein instabiler Fokus  
( $\text{Re } \lambda > 0, \text{Im } \lambda \neq 0$  für  $b \neq 0$ )

$$\text{2. Fall } (x_B^*, y_B^*) = (b, +\sqrt{1-b^2})$$

$$\underline{J}_B = \begin{pmatrix} -2b^2 + \sqrt{1-b^2} & -2b\sqrt{1-b^2} \\ -2b\sqrt{1-b^2} - b & -2 + 2b^2 \end{pmatrix}$$

Eigenwerte als Lösung der charakteristischen Gleichung:

$$\det(\underline{J}_B - \lambda \underline{1}) = \dots = (\lambda + 2)(\lambda - \sqrt{1-b^2}) = 0$$

$$\text{Eigenwerte } \lambda_1 = -2, \lambda_2 = \sqrt{1-b^2} > 0 \text{ für } |b| < 1$$

$\Rightarrow$  1 stabile, 1 instabile Richtung  $\Rightarrow$  Sattelpunkt

$$\text{3. Fall: } (x_C^*, y_C^*) = (b, -\sqrt{1-b^2})$$

$$\det(\underline{J}_C - \lambda \underline{1}) = \dots = (\lambda + 2)(\lambda + \sqrt{1-b^2}) = 0$$

$$\Rightarrow \text{Eigenwerte } \lambda_1 = -2, \lambda_2 = -\sqrt{1-b^2} < 0 \text{ für } |b| < 1$$

$\Rightarrow$  2 stabile Richtungen  $\Rightarrow$  Knoten

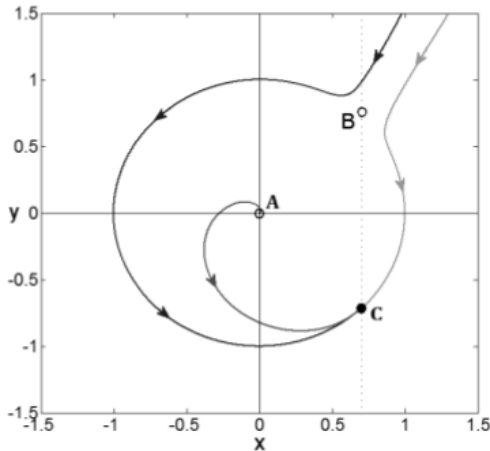
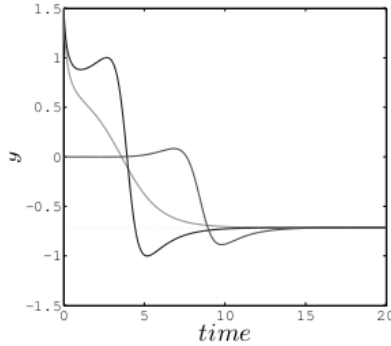
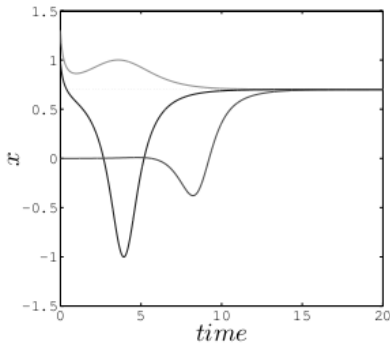
$$\text{Bei } |b|=1: (x_B^*, y_B^*) = (1, 0) = (x_C^*, y_C^*)$$

$\Rightarrow$  Fixpunkte kollidieren mit  $\text{Im } \lambda = 0$

$\Rightarrow$  intrinsische  $T \sim \frac{1}{\text{Im } \lambda} \rightarrow \infty$   
Zeit

Dynamische Szenarien:

(i) unterhalb der Bifurkation  $|b| < 1$ , z.B.  $b = 0.7$



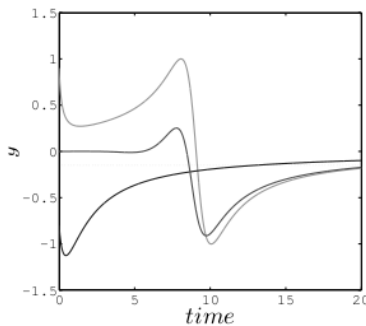
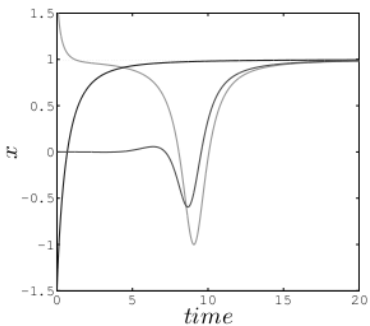
3 Fixpunkte

Fokus A

Sattelpunkt B

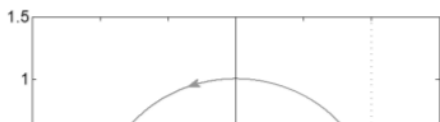
Stabiler  
Knoten C

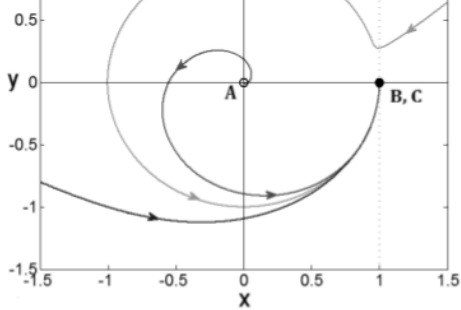
(ii) an der Bifurkation  $|b| = 1$ , z.B.  $b = 1$



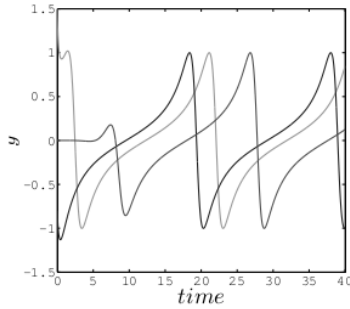
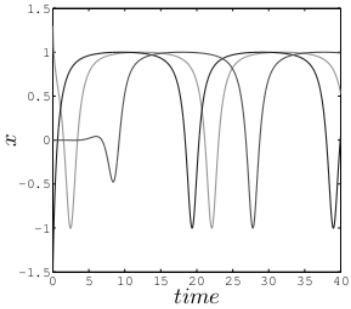
A; Fokus

Fixpunkte B & C kollidieren





(iii) oberhalb der Bifurkation  $|b| > 1$ , z.B.  $b = 1.05$

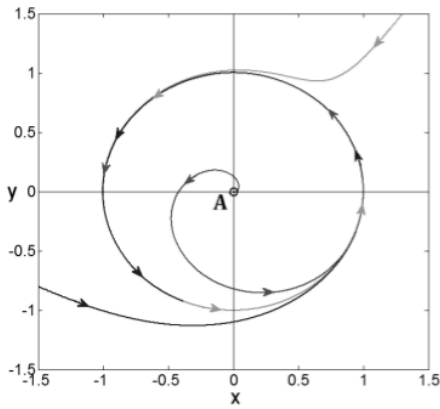


Fokus A

Grenzyklus mit Radius 1

langsame Dynamik in der

Nähe von  $(1,0) \Rightarrow$  "Geist"



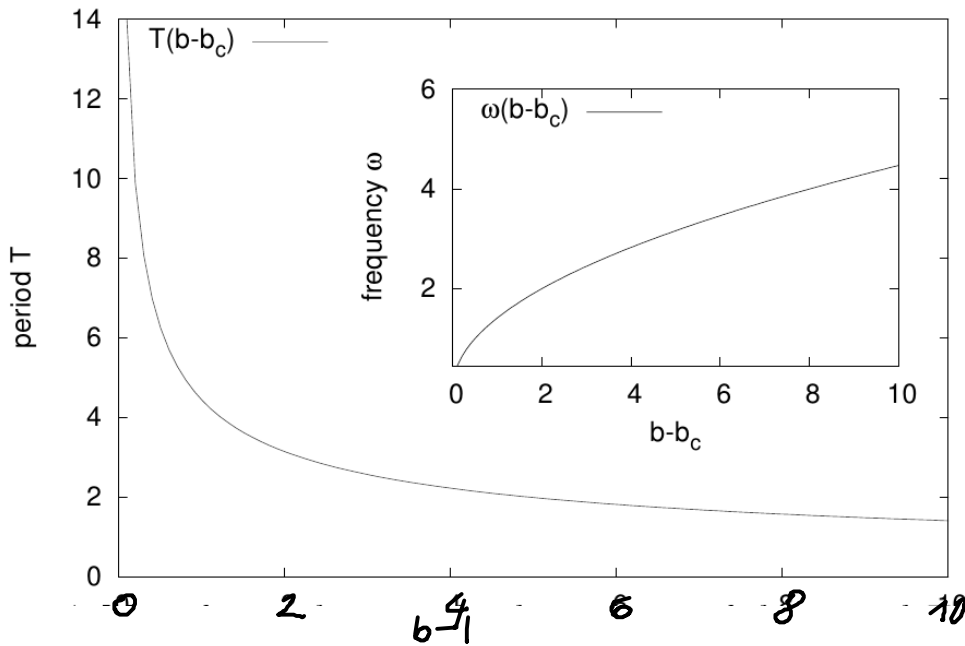
Berechnung der Periode des Grenzyklus ( $b > 1$ )

$$\dot{\varphi} = b - \cos \varphi \Rightarrow d\varphi = (b - \cos \varphi) dt$$

$$\Rightarrow dt = \frac{d\varphi}{b - \cos \varphi}$$

$$\Rightarrow \text{Periode } T = \int_0^{2\pi} \frac{d\varphi}{b - \cos \varphi} = \dots = \frac{2\pi}{\sqrt{b^2 - 1}}$$

$\Rightarrow$  Periode  $T$  divergiert für  $b \rightarrow 1$



Am Bifurkationspunkt:  
 unendliche Periode  
 endliche Amplitude  
 $\Rightarrow$  Anregbarkeit Typ I

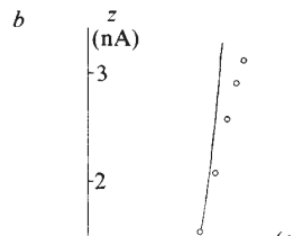
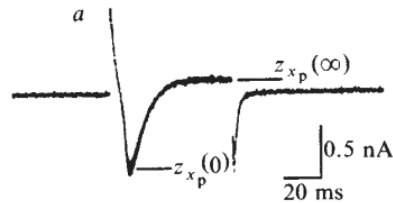
### 2.3 Hindmarsh-Rose-Modell (2D)

$$\dot{x} = c \left( x - \frac{x^3}{3} - y + z \right)$$

Kubische und quadratische Nullklare

$$\dot{y} = \frac{1}{c} (x^2 + dx - by + a)$$

- 7. Egan, R. W., Galen, P. H., Beveridge, G. C., Phillips, G. B. & Marnett, L. J. *Prostaglandins* **16**, 861-869 (1978).
- 8. Mason, R. T. & Staszewska-Barczak, J. J. *clin. exp. Pharmac. Physiol.* **6**, 678-685 (1979).
- 9. Takeguchi, C., Kohno, E. & Sih, C. J. *Biochemistry* **10**, 2372-2376 (1971).
- 10. Winter, C. A., Risley, E. A. & Nuss, G. W. *Proc. Soc. exp. Biol. Med.* **111**, 554-557 (1962).
- 11. Branceni, D., Azadian-Boulanger, G. & Jequier, R. *Archs int. Pharmacodyn.* **152**, 15-24 (1964).
- 12. Kemper, F. & Ameln, G. *Z. ges. exp. Med.* **131**, 407-415 (1959).
- 13. Flückiger, E., Schlach, W. & Taeschler, M. *Schweiz. med. Wschr.* **93**, 1232-1237 (1963).



### **A model of the nerve impulse using two first-order differential equations**

**J. L. Hindmarsh & R. M. Rose**

Department of Applied Mathematics and Astronomy and  
 Department of Physiology, University College, Cardiff,  
 Cardiff CF1 1XL, UK

$$\dot{x} = -a(f(x) - y - z) \tag{7}$$

$$\dot{y} = b(f(x) - q e^{rx} + s - y) \tag{8}$$

where  $f(x) = cx^3 + dx^2 + ex + h$ , and  $a-h, q, r$  and  $s$  are constants.

After measuring  $a$  and  $b$ , and fitting cubic and exponential functions to the  $z_{sp}(0)$  and  $z_{sp}(\infty)$  data of Fig. 1b, the solutions of equations (7) and (8) were obtained by numerical integration.

## BIFURCATIONS IN TWO-DIMENSIONAL HINDMARSH-ROSE TYPE MODEL

SHIGEKI TSUJI\*

<sup>†</sup>Aihara Complexity Modelling Project, ERATO,  
JST, 3-23-5-201 Uehara, Shibuya-ku,  
Tokyo 151-0064, Japan

TETSUSHI UETA

Center for Advanced Information Technology,  
The University of Tokushima, 2-1 Minami-Josanjima,  
Tokushima 770-8506, Japan

HIROSHI KAWAKAMI

The University of Tokushima, 2-24, Shinkura,  
Tokushima 770-8501, Japan

HIROSHI FUJII

Department of Information and Communication Sciences,  
Kyoto Sangyo University, Kamigamo-Motoyama,  
Kita-ku, Kyoto 603-8555, Japan

KAZUYUKI AIHARA<sup>†</sup>

<sup>\*</sup>Institute of Industrial Science, The University of Tokyo,  
4-6-1 Komaba, Meguro-ku, Tokyo 153-8505, Japan

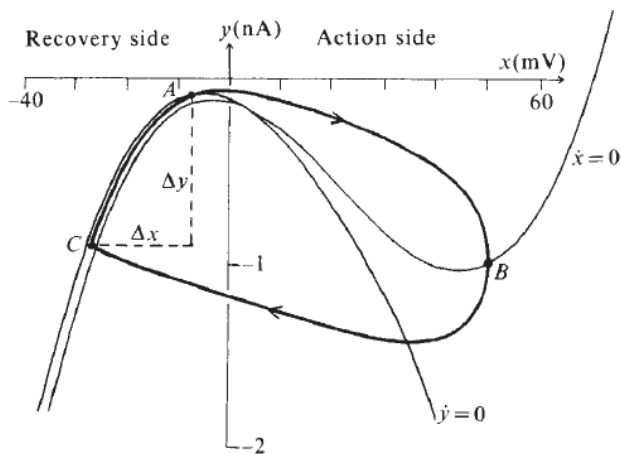


Fig. 3 Phase plane representation of the limit cycle solution to equations (7) and (8). The values of the constants are the same

Bestimmung der FP :  $\dot{x} = 0, \dot{y} = 0$

$\Rightarrow$  Lösen der Gleichung :  $\alpha x^3 + \beta x^2 + \gamma x + \delta = 0$

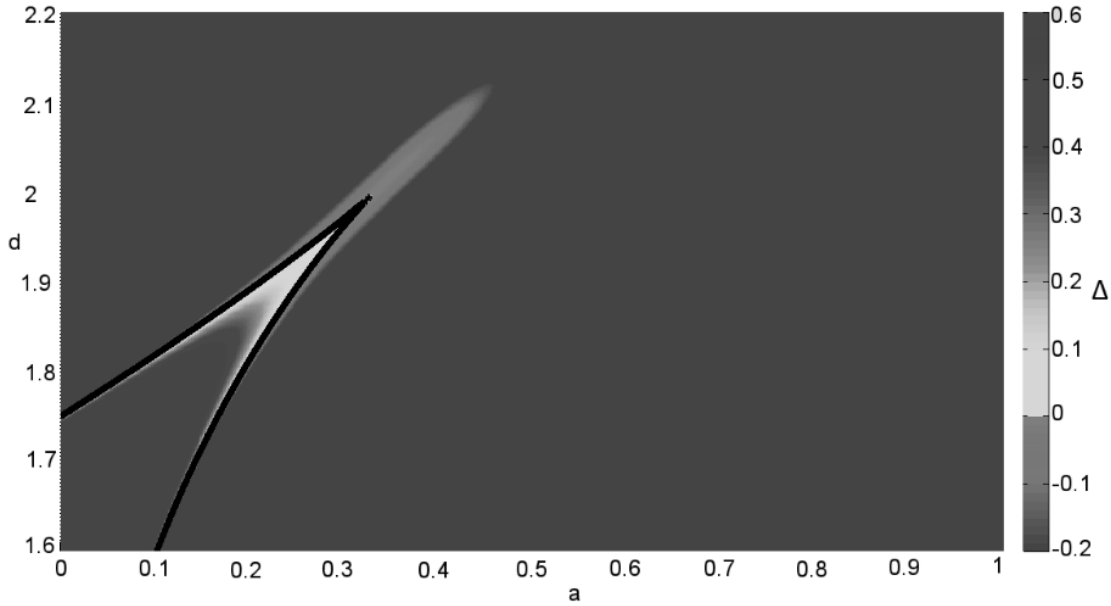
hier (für  $b=1, c=3, z=0$ ) :  $\alpha = 1, \beta = 3, \gamma = 3(d-1) + 3a$

$\Rightarrow$  Berechne Diskriminante:  $\Delta = \beta^2 \gamma^2 - 4\alpha \gamma^3 - 4\beta^3 \delta - 27\alpha^2 \delta^2 + 18\alpha \beta \gamma \delta$

(a)  $\Delta > 0$  : 3 reelle Lösungen

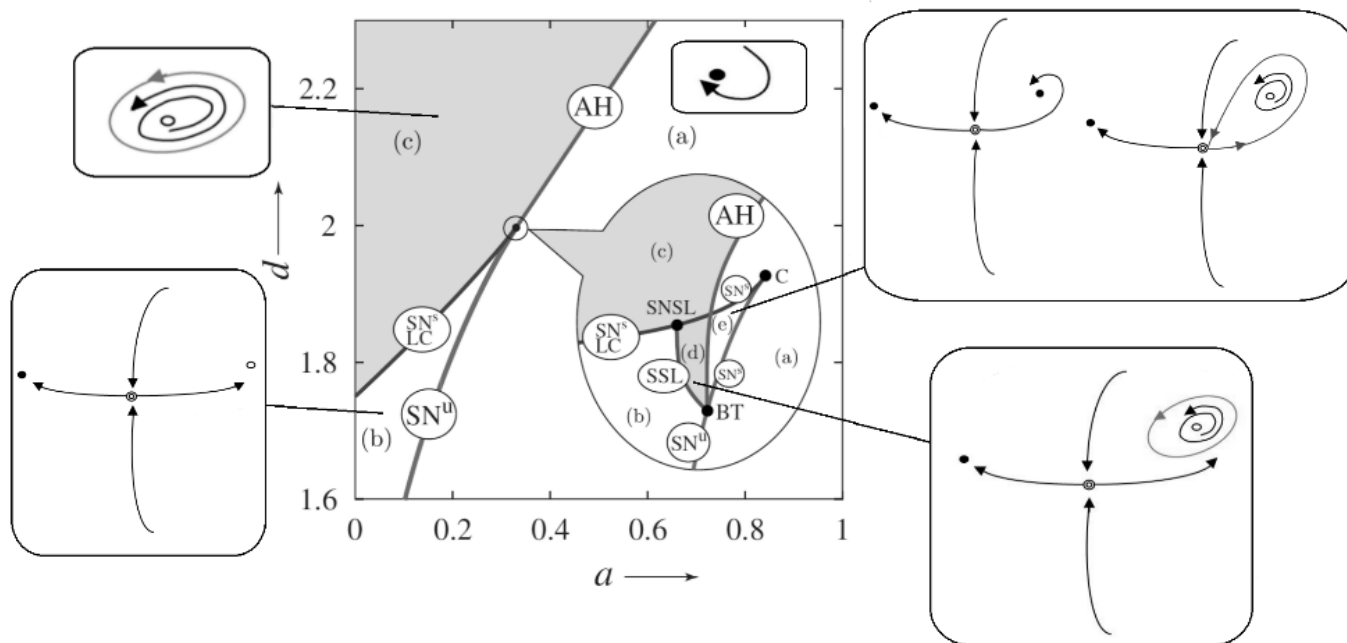
(b)  $\Delta = 0$  : 1 Lösung mit Vielfachheit 2, 1 weitere Lösung

(c)  $\Delta < 0$  : 1 reelle Lösung, 2 komplex konjugierte Lösungen



Eigenwerte per linearer Stabilitätsanalyse

$$J = \begin{pmatrix} c - c(x^*)^2 & -c \\ \frac{1}{c}(2x^* + d) & -\frac{b}{c} \end{pmatrix}$$





Label	Description
(a)	1 stable fixed point
(b)	1 saddle, 1 stable, 1 unstable fixed point
(c)	1 unstable fixed point, 1 stable limit cycle
(d)	1 saddle, 1 stable fixed points, 1 unstable fixed point, 1 stable limit cycle
(e)	1 saddle, 2 stable fixed points
AH	(Andronov-)Hopf bifurcation
SNLC	Saddle-node bifurcation on a limit cycle
SN	Saddle-node bifurcation (of equilibria)
BT	Bogdanov-Takens bifurcation
C	Cusp bifurcation
SSN	Saddle-separatrix loop bifurcation
SNSL	Saddle-node on separatrix loop bifurcation

Fortsetzung folgt...