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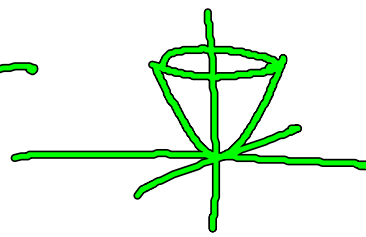
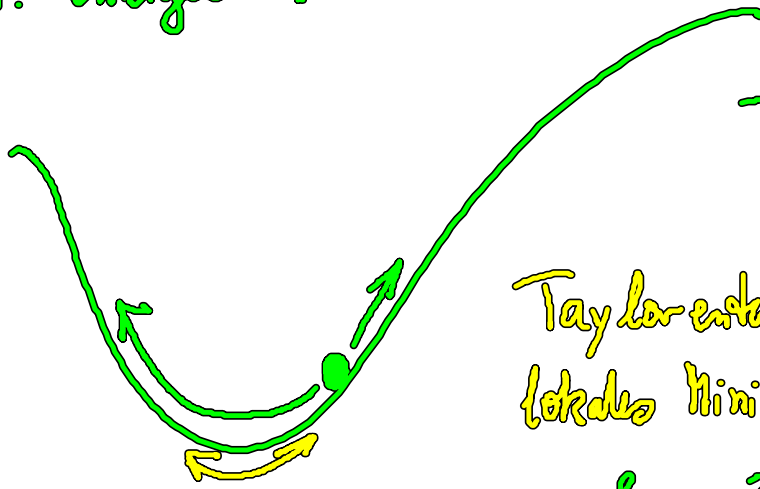
Kleine Schwingungen

Lagrange, verallg. K0



x_1, \dots, x_2

pot. Energie V



Taylorreihe von $V(\underline{x})$ um
lokales Minimum \underline{x}_0 ,

$$V(\underline{x}) = V(\underline{x}_0) + \frac{1}{2!} \sum_{i,j=1}^f \frac{\partial^2 V}{\partial x_i \partial x_j} \Big|_{\underline{x}_0} (x_i - x_i^{(0)}) (x_j - x_j^{(0)}) \\ + \dots \quad \underline{q} = \underline{x} - \underline{x}_0, \quad \underline{x} = (x_1, \dots, x_f)^T$$

Nur nächstes

$$V(\underline{x}) \rightarrow \underbrace{V(\underline{x}_0)}_{\text{Konstante}} + \frac{1}{2} \underline{q}^T \underline{V} \underline{q}$$

Konstante, wird weggelassen.

potentielle Energie

$$V(\underline{q}) = \frac{1}{2} \underline{q}^T \underline{V} \underline{q}$$

$f \times f \rightarrow$

kin. Energie hat Form

$$T(\underline{q}, \dot{\underline{q}}) = \frac{1}{2} \dot{\underline{q}}^T \underline{T}(\underline{q}) \dot{\underline{q}}$$

z.B.

$$T = \frac{1}{2} (\dot{q}_1, \dot{q}_2) \begin{matrix} f \times f \\ \begin{pmatrix} m_1 & \\ & m_2 \end{pmatrix} \end{matrix} \begin{pmatrix} \dot{q}_1 \\ \dot{q}_2 \end{pmatrix}$$
$$= \frac{1}{2} \dot{q}_1^2 m_1 + \frac{1}{2} \dot{q}_2^2 m_2.$$

⇒ Lagrange-Funktion

$$L = T - V = \frac{1}{2} \dot{\underline{q}}^T \underline{T} \dot{\underline{q}} - \frac{1}{2} \underline{q}^T \underline{V} \underline{q}$$

Normalmoden

linear Trip $\underline{q} = A \underline{Q}$, A $f \times f$ Matrix (regulär)

$$A = (\underline{a}_1, \underline{a}_2, \dots, \underline{a}_f)$$

Spaltenvektoren \underline{a}_i : Vektoren der Normalmoden

z.B. $\underline{Q} = \begin{pmatrix} 1 \\ 0 \\ \vdots \\ 0 \end{pmatrix} \rightarrow \underline{q} = \underline{a}_1.$

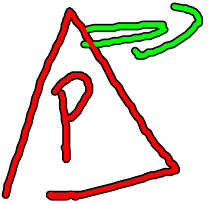
Subst matrix A so wählen, dass die Lagrangefunktion entkoppelt:

$$\begin{aligned}
 L &= \frac{1}{2} \dot{q}^T T \dot{q} - \frac{1}{2} q^T V q \\
 &= \frac{1}{2} \dot{Q}^T \dot{Q} - \frac{1}{2} Q^T \underline{D} Q \\
 &= \frac{1}{2} \sum_{i=1}^f \left(\dot{Q}_i^2 - \lambda_i Q_i^2 \right) = \sum_{i=1}^f L_i
 \end{aligned}$$

(Diagonalmatrix)

EL: $\frac{d}{dt} \frac{\partial L}{\partial \dot{Q}_i} - \frac{\partial L}{\partial Q_i} = 0; \quad L_i = \frac{1}{2} (\dot{Q}_i^2 - \lambda_i Q_i^2)$

$$\ddot{Q}_i(t) + \lambda_i Q_i(t) = 0$$



$$Q_i(t) = d_i e^{i\omega_i t} + \beta_i e^{-i\omega_i t}$$

DGL linear
harmonischer Oszillators

$$\omega_i^2 = \lambda_i$$

$$= a_i \cos \omega_i t + b_i \sin \omega_i t$$

$$L = \frac{1}{2} \dot{q}^T T \dot{q} - \frac{1}{2} q^T V q = \frac{1}{2} \dot{Q}^T \dot{Q} - \frac{1}{2} Q^T D Q$$

$$q = A Q, \quad q^T = Q^T A^T$$

$$\Rightarrow L = \frac{1}{2} \dot{Q}^T \underbrace{A^T A}_{\text{matrix}} \dot{Q} - \frac{1}{2} \underbrace{Q^T A^T V A}_{\text{matrix}} Q$$

$$\Rightarrow \begin{cases} A^T A = E & / \text{ D ist Einheitsmatrix} \\ \underline{A^T V A} = \underline{D} & \text{ D ist Diagonalmatrix.} \end{cases}$$

Simultane Eigenwertprobleme:

$$D = A^T \underline{V A} = A^T \underline{I A D}$$

$$\underline{V A} = \underline{I A D}, \quad \underline{A} = (\underline{a}_1, \dots, \underline{a}_f)$$

$$\underline{V}(\underline{a}_1, \dots, \underline{a}_f) = T(\underline{a}_1, \dots, \underline{a}_f) \underline{D}$$

$$= T(\lambda_1 \underline{a}_1, \lambda_2 \underline{a}_2, \dots, \lambda_f \underline{a}_f)$$

$$(\underline{a}_1, \underline{a}_2)_D = \left(\begin{pmatrix} a_1 \\ a_2 \end{pmatrix} \right)_D = \left(\begin{pmatrix} a_1 \\ a_2 \end{pmatrix} \right) \begin{pmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{pmatrix}$$

Damit $\underline{V} \underline{a}_i = \lambda_i \underline{I} \underline{a}_i$ Vollstg.
Eigenwertgleichung.

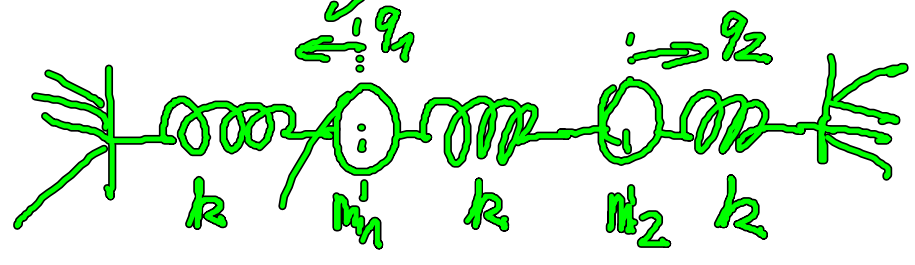
λ_i : Eigenwerte

Für nicht-triviale Lösungen der linearen Gleichungen

$(\underline{V} - \lambda_i \underline{I}) \underline{a}_i = 0$ muß die Determinante

$\det(\underline{V} - \lambda_i \underline{I}) = 0.$

Skiz.
Beispiel:



$L = \frac{1}{2} m_1 \dot{q}_1^2 + \frac{1}{2} m_2 \dot{q}_2^2 - \left[\frac{1}{2} k q_1^2 + \frac{1}{2} k q_2^2 \right.$

$\left. + \frac{1}{2} k (q_1 - q_2)^2 \right] - 2 q_1 q_2 =$

also $\underline{I} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix}$

kin. Energie - Matrix
 $= -q_1 q_2 - q_2 q_1$

$\underline{V} = k \begin{pmatrix} 2 & -1 \\ -1 & 2 \end{pmatrix} \begin{pmatrix} 2q_1 - q_2 \\ -q_1 + 2q_2 \end{pmatrix}$

$$[] = \frac{1}{2} k \begin{pmatrix} q_1 & q_2 \end{pmatrix} \begin{pmatrix} 2 & -1 \\ -1 & 2 \end{pmatrix} \begin{pmatrix} q_1 \\ q_2 \end{pmatrix} \quad (m_1 = m_2 = m)$$

$$\det(V - \lambda T) = 0 \Rightarrow \begin{vmatrix} 2k - m\lambda & -k \\ -k & 2k - m\lambda \end{vmatrix} = 0$$

$$\Rightarrow (2k - m\lambda)^2 - k^2 = 0$$

$$\Rightarrow -2k + m\lambda = \pm k$$

$$\lambda_{1,2} = \frac{k}{m}, \quad 3 \frac{k}{m}$$

$$\omega_i^2 = \lambda_i \Rightarrow \omega_1 = \sqrt{\frac{k}{m}}; \quad \omega_2 = \sqrt{3 \frac{k}{m}}$$

Eigenvektoren

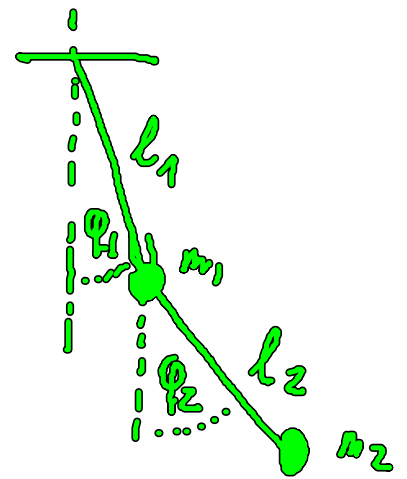
$$\begin{aligned} (V - \lambda_1 T) \underline{a}_1 = 0 &\Rightarrow \underline{a}_1 = \begin{pmatrix} 1 \\ 1 \end{pmatrix} \text{ zu } \underline{v}_1 \\ (V - \lambda_2 T) \underline{a}_2 = 0 &\Rightarrow \underline{a}_2 = \begin{pmatrix} 1 \\ -1 \end{pmatrix} \end{aligned}$$

Weiteres Beispiel: Ebene's Doppelpendel

Winkel $(q_1, q_2) = (\varphi_1, \varphi_2)$

$$T = \begin{pmatrix} (m_1 + m_2) l_1^2 & m_2 l_1 l_2 \\ m_2 l_1 l_2 & m_2 l_2^2 \end{pmatrix}$$

$$V = \begin{pmatrix} (m_1 + m_2) g l_1 & 0 \\ 0 & m_2 g l_2 \end{pmatrix}$$



Normalformen für Hamiltonfunktionen

$$L = \frac{1}{2} \dot{Q}^T \dot{Q} - \frac{1}{2} Q^T \underline{\underline{D}} Q, \quad \underline{\underline{D}} = \begin{pmatrix} \omega_1^2 & & & \\ & \omega_2^2 & & \\ & & \ddots & \\ & & & \omega_n^2 \end{pmatrix}$$

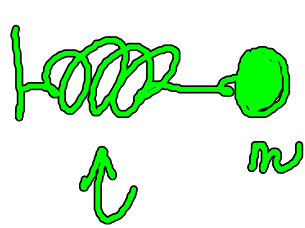
kan. Impulse $p_i = \frac{\partial L}{\partial \dot{Q}_i} = \dot{Q}_i$

$$\Rightarrow H = \sum_{i=1}^f \left(\frac{1}{2} p_i^2 + \frac{1}{2} \omega_i^2 Q_i^2 \right)$$

$$= \sum_{i=1}^f H_i; \quad H_i : i\text{-te harmonischer Oszillator}$$

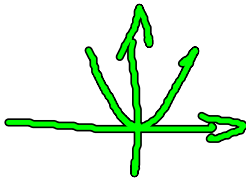
alle entkoppelt.

Ham. Oszillator



$$m\ddot{x} + m\omega^2 x = m f(t)$$

äußer Kraft

Massen im Potentialsystem  $\ddot{x} + \omega^2 x = f(t)$

$$V(x, t) = \frac{1}{2} m \omega^2 x^2 - x \cdot m f(t)$$

$$H(t) = \frac{p^2}{2m} + V(x, t)$$

$$\Rightarrow \dot{x} = \frac{\partial H}{\partial p} = \frac{p}{m}$$

$$\dot{p} = -\frac{\partial H}{\partial x} = -m\omega^2 x + m f(t)$$

$$\ddot{x} = \frac{\dot{p}}{m} = -\omega^2 x + f(t)$$

Reibungssystem:

$$\ddot{x} + \frac{1}{\tau} \dot{x} + \omega_0^2 x = f(t), \quad \tau > 0$$

Dämpfungsterm

Mo DGL-System (1. Ordnung).

$$y'(t) = \underline{A} y(t) + \underline{b}(t)$$

$$y(t) = \begin{pmatrix} x(t) \\ p(t)/m \end{pmatrix}; \quad \underline{b}(t) = \begin{pmatrix} 0 \\ f(t) \end{pmatrix}$$

$$\underline{A} = \begin{pmatrix} 0 & 1 \\ -\omega_0^2 & -1/c \end{pmatrix}$$

Homogener Fall: Zeitentwicklung.

$$(*) \quad y'(t) = A(t)y(t); \quad y = (y_1, \dots, y_n)$$

Wie beim Hamiltonschen Fluss $\phi^t: y(0) \rightarrow y(t)$

$$y(t_0) \rightarrow y(t) = \underline{U}(t, t_0) y(t_0)$$

↑
Anfangswert.

reelle $n \times n$ -Matrix.

Einsetzen in $*$: $\frac{d}{dt} (\underline{U}(t, t_0) y(t_0)) - A(t) \underline{U}(t, t_0) y(t_0)$

$$= \left(\frac{d}{dt} \underline{U}(t, t_0) - A \underline{U}(t, t_0) \right) y(t_0) = 0$$

0, es folgt

$$y(t_0) \rightarrow y(t_0)$$

$$\frac{d}{dt} \underline{U}(t, t_0) = \underline{A}(t) \underline{U}(t, t_0)$$

$U(t_0, t_0) = E$ Einheitsmatrix

$U(t, t_0)$: Zeitentwicklungsoperator